2D Spinodal Decomposition in Forced Turbulence

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What is Spinodal Decomposition?



Cahn-Hilliard Equation

• The order parameter: the concentration:

$$\psi(\mathbf{r},t) \equiv [\rho_A(\mathbf{r},t) - \rho_B(\mathbf{r},t)]/\rho$$

• The Landau-Ginzburg Free Energy:

$$\Phi(\psi) = \int d\mathbf{r} \left(-\frac{1}{2}\psi^2 + \frac{1}{4}\psi^4 + \frac{\xi^2}{2}|\nabla\psi|^2\right)$$

• Fick's Law:

$$\bar{J} = -D
abla \mu_{\rm c}$$

Chemical Potential:

$$\mu = \frac{\delta \Phi}{\delta \psi} = -\psi + \psi^3 - \xi^2 \nabla^2 \psi$$



Cahn-Hilliard Equation

• Cahn-Hilliard Equation:

$$\partial_t \psi + \mathbf{v} \cdot \nabla \psi = D \nabla^2 (-\psi + \psi^3 - \xi^2 \nabla^2 \psi)$$

• Modified Navier-Stokes Equation:

$$\partial_t \mathbf{v} + \mathbf{v} \cdot \nabla \mathbf{v} = -\frac{1}{\rho} \nabla p - \frac{\xi^2}{\rho} \nabla^2 \psi \nabla \psi + \nu \nabla^2 \mathbf{v}$$

Spinodal Decomposition vs 2D MHD

Incompressible 2D spinodal decomposition:

$$\partial_t \psi + (\nabla \phi \times \hat{\mathbf{z}}) \cdot \nabla \psi = D \nabla^2 (-\psi + \psi^3 - \xi^2 \nabla^2 \psi)$$
$$\frac{\partial}{\partial t} \nabla^2 \phi + (\nabla \phi \times \hat{\mathbf{z}}) \cdot \nabla \nabla^2 \phi = \frac{\xi^2}{\rho} (\nabla \psi \times \hat{\mathbf{z}}) \cdot \nabla \nabla^2 \psi + \nu \nabla^2 \nabla^2 \phi$$

• Incompressible 2D MHD:

$$\partial_t A + (\nabla \phi \times \hat{\mathbf{z}}) \cdot \nabla A = \eta \nabla^2 A$$
$$\frac{\partial}{\partial t} \nabla^2 \phi + (\nabla \phi \times \hat{\mathbf{z}}) \cdot \nabla \nabla^2 \phi = \frac{1}{\mu_0 \rho} (\nabla A \times \hat{\mathbf{z}}) \cdot \nabla \nabla^2 A + \nu \nabla^2 \nabla^2 \phi$$

Spinodal Decomposition vs 2D MHD

• Similarities:

Spinodal Decomposition	2D MHD
ψ	A
ξ^2	$\frac{1}{\mu_0}$
D	η^{μ_0}

- Differences:
 - Three terms in chemical potential
 - $\psi \in [-1, 1]$ while A doesn't have such restriction
- So we can use theories of 2D MHD to study spinodal decomposition!

Conserved Quantities

- Energy: $E = \int (\frac{\mathbf{v}^2}{2} + \frac{\xi^2 \mathbf{B}_{\psi}^2}{2}) \, \mathrm{d}^2 x = \int [\frac{(\nabla \phi)^2}{2} + \frac{\xi^2 (\nabla \psi)^2}{2}] \, \mathrm{d}^2 x$
- Cross Helicity: $H_C = \int \mathbf{v} \cdot \mathbf{B}_{\psi} d^2 x = \int \nabla \phi \cdot \nabla \psi d^2 x$
- Mean Square Concentration: $C_{\psi} = \int \psi^2 d^2x$
- One to one corresponding to 2D MHD conserved quantities: Energy, cross helicity, and mean square magnetic potential.
- Since negative terms in the equations, the conserved quantities do not always decay.

$$\begin{aligned} \frac{dE}{dt} &= \int (\mathbf{B}_{\psi} \cdot \xi^2 D \nabla^2 (-\mathbf{B}_{\psi} + 3\psi^2 \mathbf{B}_{\psi} - \xi^2 \nabla^2 \mathbf{B}_{\psi})) \, \mathrm{d}^2 x - \nu \int \omega^2 \, \mathrm{d}^2 x \\ &= -D \int (-\xi^{-2} j_{\psi}^2 - 6\psi B_{\psi}^2 j_{\psi} + 3\xi^{-2} \psi^2 j_{\psi}^2 + (\nabla j_{\psi})^2) \, \mathrm{d}^2 x - \nu \int \omega^2 \, \mathrm{d}^2 x \\ \frac{dH_C}{dt} &= -\nu \int \xi^{-2} j_{\psi} \omega \, \mathrm{d}^2 x - D \int (-\xi^{-2} j_{\psi} + 3\psi^2 \xi^{-2} j_{\psi} - 6\psi B_{\psi}^2 - \nabla^2 j_{\psi}) \omega \, \mathrm{d}^2 x \\ \frac{dC_{\psi}}{dt} &= \int 2\psi D \nabla^2 (-\psi + \psi^3 - \xi^2 \nabla^2 \psi) \, \mathrm{d}^2 x \\ &= -2D \int [-B_{\psi}^2 + 6\psi^2 B_{\psi}^2 + \xi^{-2} j_{\psi}^2] \, \mathrm{d}^2 x \end{aligned}$$



Spectra for 2D MHD

• Double cascade. (Biskamp 2003)



Spectra for 2D MHD

• Direct energy cascade: Alfven effect:

$$E_k = C_{IK} (\epsilon v_A)^{1/2} k^{-3/2}$$

$$A_k^2 = k^{-2} E_k \sim \epsilon^{1/2} k^{-7/2}$$

• Inverse A² cascade:

$$A_k^2 \sim \epsilon_A^{2/3} k^{-7/3}$$

Spectra for Spinodal Decomposition



Characteristic Length Scale

Define structure function:

 $S_k(t) = \langle |\psi_{\mathbf{k}}(t)|^2 \rangle$

Then the characteristic length scale is defined by:

$$L(t) = \left(\frac{\int kS_k(t) \,\mathrm{d}^2 x}{\int S_k(t) \,\mathrm{d}^2 x}\right)^{-1}$$



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Characteristic Length Scale



 L~t^{2/3} can be derived by balancing inertial term and interfacial force term.

Hinze Scale

- If external forcing at large scale is present, a direct cascade will exist. This can arrest the growth of characteristic length scale.
- The final length scale will be stabilized at Hinze Scale, by balancing the turbulent kinetic energy (break up large bubbles) and surface tension energy (stick small bubbles together).

$$L_H \sim \left(\frac{\rho}{\sigma}\right)^{-3/5} \epsilon^{-2/5}$$

- The competition of these two should have effects on spectra and cascades.
- Need to investigate further by simulation and theory.

Intermittency





Thanks!